

# Hyperfine Structure in $\text{Na}_2^+$ via Electron Spin Resonance at 4 K<sup>†</sup>

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The  $\text{Na}_2^+$  ion was produced in neon plasma and isolated in solid neon at 4 K. Its X-band electron-spin-resonance (ESR) spectrum was analyzed to yield the isotropic hyperfine interaction of 441 MHz with the  $^{23}\text{Na}$  ( $I = 3/2$ ) nucleus. An approximate interpretation of this value points to considerable  $p\sigma$  character in the predominantly  $s\sigma$  wave function of the single valence electron in the ground  $^2\Sigma_g^+$  state.

## Introduction

The one-valence-electron molecule  $\text{Na}_2^+$  has been the continuing subject of much theoretical<sup>1–15</sup> and experimental<sup>16–21</sup> study. It has been modeled<sup>1,15</sup> quite successfully as one electron in the field of two polarizable  $\text{Na}^+$  cores. Efforts have been made to establish potential energy curves over a large range of internuclear distances through ab initio and pseudopotential calculations because of the increasing accuracy of the experimental data. Presently, its vibrational and internuclear distance parameters in the ground  $^2\Sigma_g^+$  state are known,<sup>17</sup> but hyperfine splittings (hfs) due to interaction with the  $\text{Na}$  ( $I = 3/2$ ) nuclei have not been observed or calculated.

Here  $\text{Na}_2^+$  has been produced by laser vaporization of sodium metal into neon plasma, similar to that described by Jacox<sup>22</sup> and by Knight.<sup>23</sup> Its electron-spin-resonance (ESR) spectrum was subsequently measured by trapping it in solid neon at 4 K.

## Experimental Section

A Bruker (ESP 300E) ESR spectrometer was used with the matrix prepared at 4 K on a copper rod, which was then lowered into the X-band cavity ( $\nu \approx 9$  GHz). The usual ESR apparatus and accompanying cryogenics have been previously described.<sup>24</sup> Here sodium metal (Aldrich 99.9% pure) was vaporized using the 1064 nm fundamental of a pulsed Q-switched Nd:YAG laser. The sodium vapor was mixed with the neon plasma, generated by a 2450 MHz microwave discharge, just prior to condensation on the cold rod. Deposition took about 30 min with the rod at 4 K.

## Results and Analysis

The ESR spectrum of  $\text{Na}_2^+$  is shown in Figure 1. Also seen are signals resulting from sodium atoms, H atoms, and some unidentifiable signals at  $g = 2$ . The lines observed span from about 2500 to 3800 G. Assuming two equivalent nuclei with  $I = 3/2$ , this allows  $J$  to take on values of 0, 1, 2, and 3. This will result in  $m_J$  ranging from  $+J$  to  $-J$  in each case. The large hyperfine interaction in  $\text{Na}_2^+$  causes splitting of lines that would normally overlap. Thus,  $J = 0$  gives one line at the  $g$  of the molecule.  $J = 1$  will give three lines which fall at  $g$  plus a second-order term that varies with  $m_J$ ;<sup>25</sup> correspondingly,  $J = 3$  will result in 7 lines. There are a total of 32 predicted lines

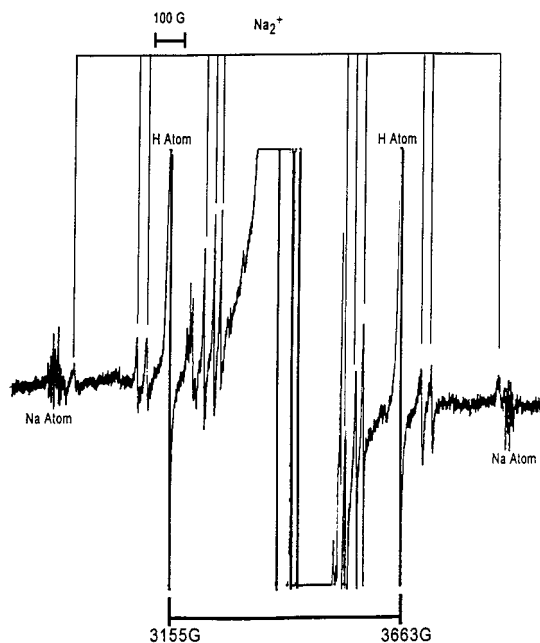


Figure 1. ESR spectrum of  $\text{Na}_2^+$  in solid neon at 4 K.

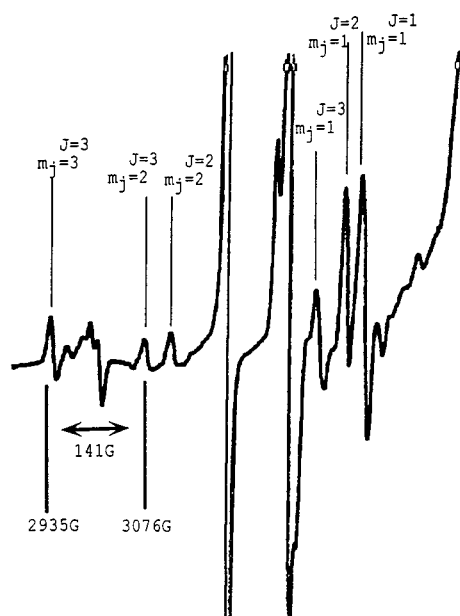
including both perpendicular and parallel contributions. However, only 23 assignments can be made due to interference at  $g = 2$  and overlap of Na atom lines at high field. Figure 2 shows an expanded view of two regions of the  $\text{Na}_2^+$  spectrum.

The perpendicular lines as indicated in Table 1 were fit by obtaining the eigenvalues generated by the appropriate spin Hamiltonian for this system, which is given in eq 1.<sup>25</sup> To

$$H = g_{\parallel}\beta H_z S_z + g_{\perp}\beta(H_x S_x + H_y S_y) + A_{\parallel}J_z S_z + A_{\perp}\beta(J_x S_x + J_y S_y) \quad (1)$$

evaluate  $H_{\perp}$  or  $H_{\parallel}$ , a matrix of the order  $(2S + 1)(2J + 1) = 2(2J + 1)$  was solved, where  $J = |I_1 + I_2|, \dots, |I_1 - I_2|$ , and  $M_J = J, J - 1, \dots, -J$ , since  $S = 1/2$ ,  $M_S = \pm 1/2$ . The spin matrix was set up individually for each  $I$  value. A suitable fit was generated by carefully varying the magnetic parameters  $A_{\perp}$  and  $A_{\parallel}$  and then comparing the calculated to the observed perpendicular line positions. The parallel parameters influenced this fit but to a lesser degree, so that they could not be reliably determined in this way. The weak parallel lines were difficult to identify and assign; however, one obtains  $g_{\parallel} = 2.0006(20)$  and  $A_{\parallel} =$

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**Figure 2.** Expanded view of the ESR spectrum of  $\text{Na}_2^+$  in solid neon, showing the low-field perpendicular lines.

**TABLE 1: Observed Perpendicular Line Positions of the ESR Spectrum of  $\text{Na}_2^+$  in Solid Neon Compared to the Calculated Values ( $\nu = 9.5855$  GHz)**

$J, M_j$	$H_{\perp}$ obs ( $\pm 1$ G)	$H_{\perp}$ calc (G) <sup>a</sup>	$J, M_j$	$H_{\perp}$ obs ( $\pm 1$ G)	$H_{\perp}$ calc (G) <sup>a</sup>
3, 3	2935.0	2935.1	1, 0	3414.3	3414.5
3, 2	3075.5	3075.7	0, 0		3421.2
2, 2	3099.9	3101.6	3, -1	3539.4	3538.8
3, 1	3222.5	3222.0	2, -1	3561.6	3561.4
2, 1	3245.0	3245.3	1, -1	3576.5	3576.8
1, 1	3259.8	3260.5	3, -2	3709.3	3709.2
3, 0	3378.0	3378.1	2, -2	3731.9	3732.6
2, 0		3400.2	3, -3		3884.0

<sup>a</sup> Calculated using the parameters given in Table 2.

**TABLE 2: Derived Magnetic Parameters for  $\text{Na}_2^+$  in Solid Neon at 4 K**

derived parameter	
$g_{\parallel}$	2.0006(20)
$g_{\perp}$	2.0017(9)
$A_{\parallel}$	155(20) G
	434(60) MHz
$A_{\perp}$	158.4(9) G
	444(3) MHz
$A_{\text{iso}}$	441(20) MHz

155(20) G, with estimated uncertainties. (Because of the large hfs, second-order perturbation<sup>26</sup> is not expected to yield accurate parameters. If applied, it gives  $A_{\perp} = 158.4$  G,  $A_{\parallel} = 164.8$  G.)

## Discussion

Values for the isotropic  $A_{\text{iso}} = (A_{\parallel} + 2A_{\perp})/3$  and dipolar  $A_{\text{dip}} = (A_{\parallel} - A_{\perp})/3$  hyperfine parameters are included in Table 2.  $A_{\text{iso}} = 441$  MHz, representing the  $s\sigma$  contribution to the valence electron wave function. When multiplied by 2 to account for the two equivalent nuclei, it is of the same magnitude as the corresponding neutral atomic value tabulated by Morton and Preston,<sup>27</sup> 927.1 MHz. In the usual approximate procedure,<sup>25</sup> one then finds  $882/927 = 95\%$   $s\sigma$  character. (These are a set of self-consistent parameters. The experimental value<sup>28</sup> for the Na atom is 885.8 MHz, which makes this calculated  $s\sigma$  character even higher.) However, here our model involves interaction with

the nuclei in two  $\text{Na}^+$  cores, where the hfs will be larger because of the positive charge. According to Goudsmit's formula,<sup>29</sup> this interaction would increase by approximately a factor of 4 from the neutral to the cationic alkali metal atom. Since the electron is shared by two nuclei here, comparison could be made with a value of  $2 \times 927.1$  MHz, yielding, probably a minimum,  $s\sigma$  character of 48%.  $A_{\text{dip}}$  is even more difficult to interpret. One can only suggest that the evidence indicates that there is appreciable  $p\sigma$  character in the unpaired-electron wave function. Because of the large uncertainty in  $A_{\parallel}$ ,  $A_{\text{dip}}$  can vary from  $-24$  to  $+18$  MHz, and therefore does not provide the usual information about the  $p\sigma$  contribution.

These considerations concerning the character of the one valence electron are similar to those considered in some detail in our earlier work on the  $\text{Cu}_2^+$ ,  $\text{Ag}_2^+$ ,  $\text{Au}_2^+$  ESR spectra.<sup>30</sup> There,  $2A_{\text{iso}}$  for  $\text{Cu}_2^+$  and  $\text{Ag}_2^+$  were 112 and 114%, respectively, of that of the neutral atom. It was estimated that the  $\sigma$  electron contained 100 and 70(20)%  $s$  character.

An interesting comparison can also be made with the  $\text{Na}_3$  molecule. The sodium trimer was observed via ESR by Lindsay et al.<sup>31</sup> and was found to have two equivalent sodium atoms in an isosceles triangle geometry. The unpaired spin density was concentrated on the equivalent atoms and was found to have predominately  $s$  character,  $A_{\text{iso}} = 420$  MHz.

Our measured  $g_{\perp}$  value (Table 2) could not be determined accurately enough to provide a reliable calculation of the mixing of excited  $\Pi$  electronic states into the ground state. The observed  $\Delta g_{\perp} = g_{\perp} - g_e$  is negative, which indicates the coupling occurs to a  ${}^2\Pi_r$  excited state, presumably the lowest observed  ${}^2\Pi_u$  state<sup>17</sup> at about  $23\,000\text{ cm}^{-1}$ . Substituting a spin-orbit coupling constant for the observed state,  $\xi = 10.9\text{ cm}^{-1}$  (or the sodium atom value =  $11.46\text{ cm}^{-1}$ ) into

$$\Delta g_{\perp} = (2\xi|\Delta E\rangle\langle\Pi|L_x|\Sigma\rangle\langle\Sigma|L_x|\Pi\rangle)$$

yields  $\Delta E = 11\,000\text{ cm}^{-1}$ .

Although there are theoretical treatments of  $\text{Li}_2^+$ ,<sup>4,9-11,13</sup> there are again no hyperfine structure data on it or the heavier alkali metal molecules.

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